Deformation Quantization in Algebraic Geometry

Amnon Yekutieli Ben Gurion University, ISRAEL

Lecture notes (with bibliography) at: http://www.math.bgu.ac.il/ ~amyekut/lectures/def-quant.html

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plan

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- 2. The Local Picture

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- 2. The Local Picture
- 3. Deforming Algebraic Varieties

1 What is Deformation Quantization?

Let \mathbb{K} be a field of characteristic 0 and let C be a commutative \mathbb{K} -algebra.

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A *Poisson bracket* on *C* is a \mathbb{K} -bilinear function

$$\{-,-\}: C \times C \to C$$

which makes *C* into a Lie algebra, and is a bi-derivation (i.e. a derivation in each argument).

Let $\mathcal{T}_C = \text{Der}_{\mathbb{K}}(C)$ be the module of derivations of *C*. Given $\alpha \in \bigwedge_C^2 \mathcal{T}_C$ we can define a bi-derivation $\{-, -\}_{\alpha}$ by

$$\{f,g\}_{\alpha} := \langle \mathrm{d}f \wedge \mathrm{d}g, \alpha \rangle$$

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If $\{-,-\}_{\alpha}$ is a Poisson bracket on *C* (i.e. the Jacobi identity holds) then we call α a *Poisson structure*.

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Example 1. Classical Hamiltonian mechanics. Here $\mathbb{K} = \mathbb{R}$, *X* is an even dimensional differentiable manifold (the phase space; often $X = T^*Y$, the cotangent bundle of a manifold *Y*) and $C := C^{\infty}(X)$, the ring of differentiable \mathbb{R} -valued functions. See [GS].

Example 2. Lie theory: let \mathfrak{g} be a finite dimensional Lie algebra. Then $C := \operatorname{Sym} \mathfrak{g}$ can be identified with $\mathcal{O}(\mathfrak{g}^*)$, the ring of algebraic functions on the dual \mathfrak{g}^* . There is an intrinsically defined Poisson bracket called the Kostant - Kirillov bracket.

Example 3. Here is an explicit formula. Take n = 2m and $C := \mathbb{K}[t_1, \ldots, t_n]$, the polynomial algebra. Then

$$\alpha := \sum_{i=1}^{m} \frac{\partial}{\partial t_i} \wedge \frac{\partial}{\partial t_{i+m}}$$

is a Poisson structure on C.

Example 4.

Let \hbar denote a central variable (the "Planck constant"). Suppose A is a flat $\mathbb{K}[[\hbar]]$ -algebra with multiplication \star , and $\psi: A/(\hbar) \xrightarrow{\simeq} C$ is a \mathbb{K} -algebra isomorphism.

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Given $f, g \in C$ choose arbitrary lifts $\tilde{f}, \tilde{g} \in A$, and define

$$\{f,g\}_{\star} := \psi \left(\frac{1}{2\hbar} (\tilde{f} \star \tilde{g} - \tilde{g} \star \tilde{f}) \right).$$

This makes sense because \hbar divides $\tilde{f} \star \tilde{g} - \tilde{g} \star \tilde{f}$ in A.

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The element $\{f, g\}_{\star} \in C$ turns out to be independent of the choice of lifts, and this is a Poisson bracket.

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Definition 5. A deformation quantization of *C* is a $\mathbb{K}[[\hbar]]$ -bilinear associative multiplication \star on the $\mathbb{K}[[\hbar]]$ -module $C[[\hbar]]$, of the form

$$f \star g = fg + \sum_{j=1}^{\infty} \beta_j(f,g)\hbar^j$$

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are bi-differential operators. The multiplication \star is called a star product.

The reason that the β_j have to be bi-differential operators is to make the star product \star local; otherwise the deformation has no geometric significance. The reason that the β_j have to be bi-differential operators is to make the star product \star local; otherwise the deformation has no geometric significance.

Definition 6. Given a Poisson structure α on the algebra C, a deformation quantization of the Poisson algebra (C, α) is a deformation quantization \star of C such that

$$\{f,g\}_{\star} = \{f,g\}_{\alpha}$$

for all $f, g \in C$.

The original idea by the physicists Flato et. al. [BFFLS] in 1978 was that in the setup of Example 1 the quantization process should model the transition from classical to quantum mechanics. The original idea by the physicists Flato et. al. [BFFLS] in 1978 was that in the setup of Example 1 the quantization process should model the transition from classical to quantum mechanics.

For a symplectic manifold X and $C = C^{\infty}(X)$ the problem was solved by De Wilde and Lacomte in 1983 [DL]. A more geometric solution was discovered by Fedosov in 1994 [Fe]. The original idea by the physicists Flato et. al. [BFFLS] in 1978 was that in the setup of Example 1 the quantization process should model the transition from classical to quantum mechanics.

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The general case, i.e. $C = C^{\infty}(X)$ for a Poisson manifold X, was solved by Kontsevich in 1997 [Ko1]. See the survey articles [Ke] and [CI].

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A gauge equivalence is a $\mathbb{K}[[\hbar]]$ -linear automorphism $\gamma: C[[\hbar]] \xrightarrow{\simeq} C[[\hbar]]$ of the form $\gamma = \mathbf{1} + \sum_{j=1}^{\infty} \gamma_j \hbar^j$, where $\mathbf{1}: C \xrightarrow{\simeq} C$ is the identity, and each $\gamma_j: C \to C$ is a differential operator. To state Kontsevich's result we need the notion of gauge equivalence.

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Two star products \star and \star' on $C[[\hbar]]$ are called gauge equivalent if

$$\gamma(f \star' g) = \gamma(f) \star \gamma(g)$$

for some gauge equivalence γ .

A formal Poisson structure on *C* is a series $\alpha = \sum_{j=1}^{\infty} \alpha_j \hbar^j$, with $\alpha_j \in \bigwedge_C^2 \mathcal{T}_C$, which is a Poisson structure on the commutative algebra $C[[\hbar]]$. A formal Poisson structure on *C* is a series $\alpha = \sum_{j=1}^{\infty} \alpha_j \hbar^j$, with $\alpha_j \in \bigwedge_C^2 \mathcal{T}_C$, which is a Poisson structure on the commutative algebra $C[[\hbar]]$.

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There is a similar notion of gauge equivalence of formal Poisson structures.

Example 7. If α_1 is a Poisson structure on *C* then $\alpha := \alpha_1 \hbar$ is a formal Poisson structure.

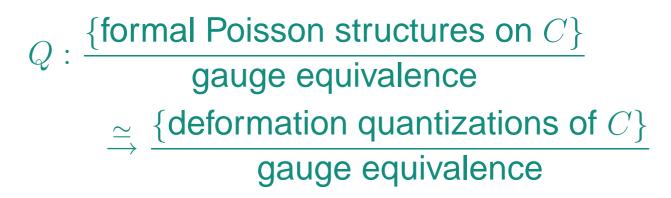
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The function Q is called the *quantization map*.

By "preserving first order terms" I mean that given a formal Poisson structure $\alpha = \sum_{j=1}^{\infty} \alpha_j \hbar^j$, and corresponding deformation quantization $\star = Q(\alpha)$, one has

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The proof has a local aspect, which I will explain next. The global part of the proof uses formal geometry, and I will talk about it later, in the context of algebraic geometry.

2 The Local Picture

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- (iii) $\mathbb{K} = \mathbb{R}$, and $C = C^{\infty}(U)$, the ring of differentiable functions on some open set $U \subset \mathbb{R}^n$ (in the classical topology), with coordinates t_1, \ldots, t_n .

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It is crucial that *C* is equipped with coordinates.

It turns out that the local deformation problem is best handled using differential graded Lie algebras. This idea is attributed to Deligne. It turns out that the local deformation problem is best handled using differential graded Lie algebras. This idea is attributed to Deligne.

A DG (differential graded) Lie algebra is a graded \mathbb{K} -module $\mathfrak{g} = \bigoplus_{i \in \mathbb{Z}} \mathfrak{g}^i$, with a bracket [-, -] satisfying the graded version of the Lie algebra identities, together with a differential $d : \mathfrak{g}^i \to \mathfrak{g}^{i+1}$ compatible with [-, -].

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Given a DG Lie algebra \mathfrak{g} , let us define a new DG Lie algebra

$$\mathfrak{g}[[\hbar]]^+ := \bigoplus_i \, \hbar \mathfrak{g}^i[[\hbar]] \subset \bigoplus_i \, \mathfrak{g}^i[[\hbar]],$$

in which \hbar is central.

The Maurer-Cartan equation in $\mathfrak{g}[[\hbar]]^+$ is

$$d(\alpha) + \frac{1}{2}[\alpha, \alpha] = 0$$

for

$$\alpha = \sum_{j=1}^{\infty} \alpha_j \hbar^j \in \mathfrak{g}^1[[\hbar]]^+.$$

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One defines

 $\mathrm{MC}(\mathfrak{g}[[\hbar]]^+) :=$

{solutions of MC equation in $\mathfrak{g}[[\hbar]]^+$ }

gauge equivalences

Let us return to our deformation problem, where C is one of the commutative \mathbb{K} -algebras (i)-(iii) above.

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Recall the module of derivations T_C . For $i \ge -1$ define

$$\mathcal{T}^{i}_{\text{poly}}(C) := \bigwedge_{C}^{i+1} \mathcal{T}_{C}.$$

So
$$\mathcal{T}_{\text{poly}}^{-1}(C) = C$$
 and $\mathcal{T}_{\text{poly}}^{0}(C) = \mathcal{T}_{C}$.

The direct sum

$$\mathcal{T}_{\text{poly}}(C) := \bigoplus_{i} \mathcal{T}^{i}_{\text{poly}}(C)$$

is a DG Lie algebra, called the algebra of poly derivations of *C*. The Lie bracket is the Schouten-Nijenhuis bracket, and the differential is 0.

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The solutions of the Maurer-Cartan equation in $\mathcal{T}_{poly}(C)[[\hbar]]^+$ are precisely the formal Poisson structures on C.

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So $\mathcal{D}_{\text{poly}}^{-1}(C) = C$ and $\mathcal{D}_{\text{poly}}^{0}(C) = \mathcal{D}(C)$, the ring of differential operators.

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$$f \star g := fg + \sum_{j=1}^{\infty} \beta_j(f,g)\hbar^j$$

is an associative deformation of the multiplication of C.

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given by

$$\mathcal{U}_1(\partial_1 \wedge \dots \wedge \partial_k)(f_1, \dots, f_k) := \frac{1}{k!} \sum_{\sigma \in S_k} \operatorname{sgn}(\sigma) \partial_{\sigma(1)}(f_1) \cdots \partial_{\sigma(k)}(f_k)$$

for $f_i \in C$ and $\partial_i \in \mathcal{T}_C$.

It is known that U_1 is a quasi-isomorphism – see [Ko1] for the case $C = C^{\infty}(U)$, and [Ye1] for the case C = O(U) – and it induces an isomorphism of graded Lie algebras in cohomology. It is known that U_1 is a quasi-isomorphism – see [Ko1] for the case $C = C^{\infty}(U)$, and [Ye1] for the case C = O(U) – and it induces an isomorphism of graded Lie algebras in cohomology.

But U_1 is not a DG Lie algebra homomorphism!

Theorem 9. (Kontsevich Formality Theorem)

Then \mathcal{U}_1 extends to an L_∞ quasi-isomorphism

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Moreover, each of the maps U_j is invariant under linear change of coordinates.

There is an induced L_∞ quasi-isomorphism

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A calculation shows that we get a bijection

$$\mathrm{MC}(\mathcal{U}) : \mathrm{MC}(\mathcal{T}_{\mathrm{poly}}(C)[[\hbar]]^{+})$$
$$\xrightarrow{\simeq} \mathrm{MC}(\mathcal{D}_{\mathrm{poly}}(C)[[\hbar]]^{+})$$

with an explicit formula.

Therefore:

Corollary 10. Assume $\mathbb{R} \subset \mathbb{K}$. In each of the cases (i) - (iii) above there is a canonical bijection of sets

 $Q: \frac{\{\text{formal Poisson structures on } C\}}{\text{gauge equivalence}}$ $\xrightarrow{\simeq} \frac{\{\text{deformation quantizations of } C\}}{\text{gauge equivalence}}$

preserving first order terms.

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A Poisson structure on X is an element $\alpha \in \Gamma(X, \bigwedge_{\mathcal{O}_X}^2 \mathcal{T}_X)$ such that the bi-derivation $\{-, -\}_{\alpha}$ is a Poisson bracket on the sheaf of functions \mathcal{O}_X .

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The aim is to deform the sheaf \mathcal{O}_X .

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such that the formula

$$f \star g := fg + \sum_{j=1}^{\infty} \beta_j(f,g)\hbar^j,$$

for local sections $f, g \in \mathcal{O}_X$, defines an associative $\mathbb{K}[[\hbar]]$ -algebra structure on the sheaf $\mathcal{O}_X[[\hbar]]$.

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But when $H^1(X, \mathcal{D}_X) = 0$, which holds in many cases, the two notions of deformation quantization coincide. Here \mathcal{D}_X is the sheaf of differential operators on *X*.

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But when $H^1(X, \mathcal{D}_X) = 0$, which holds in many cases, the two notions of deformation quantization coincide. Here \mathcal{D}_X is the sheaf of differential operators on X.

Recall that X is called \mathcal{D} -affine if $\mathrm{H}^{i}(X, \mathcal{M}) = 0$ for any quasi-coherent left \mathcal{D}_{X} -module \mathcal{M} and any i > 0.

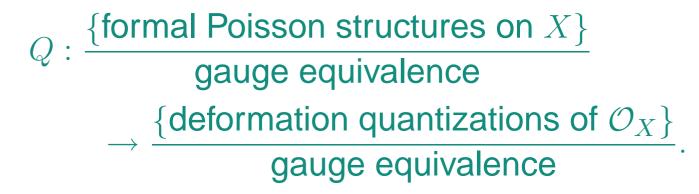
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Theorem 12. (Y) Assume X is \mathcal{D} -affine and $\mathbb{R} \subset \mathbb{K}$. Then there is a canonical function



The function Q preserves first order terms, and respects étale morphisms $X' \rightarrow X$. If X is affine then Q is bijective. There is an explicit formula for Q. Theorem 12 applies to two large classes of varieties.

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1. Affine varieties. Note that if $X = \operatorname{Spec} C$ is affine, but does not admit an étale morphism to $A^n_{\mathbb{K}}$, then this result is not trivial, since there has to be some gluing. Cf. case (i) earlier.

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- 1. Affine varieties. Note that if $X = \operatorname{Spec} C$ is affine, but does not admit an étale morphism to $A^n_{\mathbb{K}}$, then this result is not trivial, since there has to be some gluing. Cf. case (i) earlier.
- 2. The flag varieties X = G/P, where *G* is a connected reductive algebraic group and *P* is a parabolic subgroup. By the Beilinson-Bernstein Theorem the variety *X* is \mathcal{D} -affine. This class of varieties includes the projective spaces $\mathbb{P}^n_{\mathbb{K}}$.

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The universal Taylor expansion is a canonical embedding of algebras

$$\pi^{-1}\mathcal{O}_X \subset \mathcal{O}_{\operatorname{Coor} X}[[t]].$$

Due to the Formality Theorem we obtain an L_∞ quasi-isomorphism

$$\mathcal{U}: \mathcal{O}_{\operatorname{Coor} X} \mathbin{\widehat{\otimes}} \mathcal{T}_{\operatorname{poly}}(\mathbb{K}[[t]]) \to \mathcal{O}_{\operatorname{Coor} X} \mathbin{\widehat{\otimes}} \mathcal{D}_{\operatorname{poly}}(\mathbb{K}[[t]])$$

of sheaves of DG Lie algebras on Coor X.

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of sheaves of DG Lie algebras on Coor X.

If we had a section $\sigma: X \to \operatorname{Coor} X$ then we could pull \mathcal{U} down to an L_{∞} quasi-isomorphism

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However usually there are no global sections of Coor X.

The GL_n -invariance in the Formality Theorem says that it suffices to find a section $\sigma : X \to LCC X$.

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In the C^{∞} context such global sections $\sigma : X \to LCC X$ do exists (because the fibers of this bundle are contractible).

The GL_n -invariance in the Formality Theorem says that it suffices to find a section $\sigma : X \to LCC X$.

In the C^{∞} context such global sections $\sigma : X \to LCC X$ do exists (because the fibers of this bundle are contractible).

But this is not the case in algebraic geometry. (Here is where our work diverges from that of Kontsevich.) So we must use a trick.

Our geometric trick is called simplicial sections. We can choose an open covering $X = \bigcup U_i$ with sections $\sigma_i : U_i \to \operatorname{LCC} X$.

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Using an averaging process for unipotent group actions [Ye3], for any i_0, \ldots, i_q we then obtain a morphism

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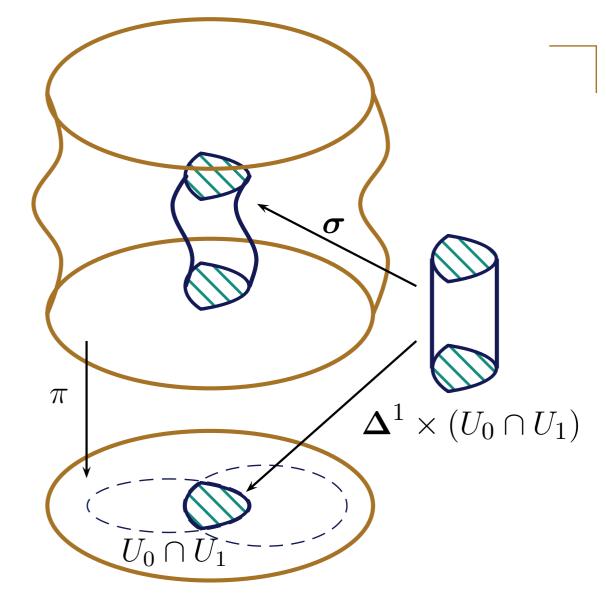
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As q varies we have a simplicial section σ , i.e. the simplicial relations are satisfied.

 $\operatorname{LCC} X$ Here is an illustration of the case q = 1. We start with sections over two open sets $\sigma_i: U_i \to \operatorname{LCC} X$ σ_0 π σ_1 U_0 X



and we pass to a simplicial section $\sigma_{(0,1)}$ interpolating between them

The simplicial section σ gives rise to an L_∞ quasi-isomorphism

 $\Psi_{\boldsymbol{\sigma}} : \operatorname{Mix}(\mathcal{T}_{\operatorname{poly},X}) \to \operatorname{Mix}(\mathcal{D}_{\operatorname{poly},X})$

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It is now pretty easy to deduce Theorem 12, including the cohomological conditions.

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- end -

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